

# Dynamic motion of La atom inside the $C_{74}$ ( $D_{3h}$ ) cage: a relativistic DFT study

Dongxu Tian · Suzhen Ren · Ce Hao

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**Abstract** The interaction between lanthanum atom (La) and  $C_{74}$  ( $D_{3h}$ ) was investigated by all-electron relativistic density function theory (DFT). With the aid of the representative patch of  $C_{74}$  ( $D_{3h}$ ), we studied the interaction between  $C_{74}$  ( $D_{3h}$ ) and La and obtained the interaction potential. Optimized structures show that there are three equivalent stable isomers, with La located about 1.7 Å off center. There is one transition state between every two stable isomers. According to the minimum energy pathway, the possible movement trajectory of La atoms in the  $C_{74}$  ( $D_{3h}$ ) cage was explored. The calculated energy barrier for La atoms moving from the stable isomer to the transition state is 18.4 kcal mol<sup>-1</sup>. In addition, the dynamic NMR spectra of La@ $C_{74}$  according to the trajectory was calculated.

**Keywords** Relativistic DFT · La@ $C_{74}$  · Energy potential surface · Dynamic NMR

## Introduction

Since the first extraction of metallofullerene La@ $C_{82}$  [1], endohedral metallofullerenes (EMFs) have attracted considerable attention [2, 3] due to their novel structures with promising properties and applications in material and biological science.  $C_{74}$  is very unstable and is called a “missing fullerene” among the fullerene family. However, the  $C_{74}$  cage is stabilized when a divalent metal (M) is inserted into the cage, with two electrons transferred from the metal to the  $C_{74}$  cage [4]. Experimentally, some divalent  $M^{2+}@C_{74}^{2-}$  EMFs (M = Ba [5–7], Eu [8], Sm [9–11], Ca [6, 12], Sr [6, 7], or Yb [13]) have been isolated.  $M^{2+}@C_{74}^{2-}$  has been characterized by UV–VIS, Raman, and EPR spectroscopy.

However, insoluble EMFs, such as La@ $C_{74}$ , have not yet been isolated, although they are observed regularly in raw soot by mass spectrometry. Nikawa et al. [3] first reported isolation of La@ $C_{74}$  as an EMF derivative, La@ $C_{74}$  ( $C_6H_3Cl_2$ ), suggesting that La@ $C_{74}$  has an electronic structure of  $La^{3+}@C_{74}^{3-}$  and is paramagnetic.

The position and motion of the encapsulated metal atom plays an important role in the chemical and physical properties of EMFs. Both experimental and theoretical studies have revealed the dynamic behavior of metal atoms encapsulated in pristine and functionalized fullerene cages. Experimentally, Miyaka et al. [14] have studied the motion of the scandium ions in  $Sc_2@C_{84}$  by <sup>45</sup>Sc solution NMR spectroscopy. The two scandium ions in  $Sc_2@C_{84}$  with  $D_{2d}$  symmetry rapidly change their positions. Nishibori et al. [15] reported that the trajectory of lanthanum atoms is like a bowl or hemisphere in  $C_{82}$  cages at room temperature using the maximum entropy method (MEM). Recently, <sup>13</sup>C NMR spectroscopy strongly suggests that two encapsulated Lu atoms rapidly rotate in  $T_d$ - $C_{76}$  fullerene cages [16]. A similar dynamical motion of encaged metal atoms in fullerenes has been observed in La<sub>2</sub>@ $C_{80}$  [17, 18]. Theoretically, Andreoni et al. [19] obtained the dynamic trajectory of La inside  $C_{60}$  by ab initio molecular dynamics. Heine et al. calculated the dynamic NMR spectra [20] of Sc<sub>3</sub>N@ $C_{80}$  using quantum Born-Oppenheimer molecular dynamics simulations, followed by DFT-NMR calculations on a large series of snapshots. Jin et al. [21] reported that the La atom probably undergoes boat-shaped movement at high temperatures. Moreover, dynamical motion of divalent metals in  $C_{74}$  ( $D_{3h}$ ) cages has been investigated. The <sup>13</sup>C NMR spectra of Ca@ $C_{74}$  indicate that the Ca atom hops inside the  $D_{3h}$  cage [22]. A similar dynamic motion has been observed in Yb@ $C_{74}$  by <sup>13</sup>C NMR spectroscopy [4]. The dynamic behavior of europium in Eu@ $C_{74}$  was investigated by quasi-relativistic density-functional based tight binding calculations and molecular dynamics simulations by Viezte et al. [23]. To the best of our knowledge, few

D. Tian · S. Ren · C. Hao (✉)  
State Key Laboratory of Fine Chemicals, Dalian University of  
Technology, Dalian 116024, People's Republic of China  
e-mail: haoce@dlut.edu.cn

theoretical studies have been reported with the purpose of exploring the dynamic behavior of lanthanum atoms in  $C_{74}$  ( $D_{3h}$ ) cages.  $La^{3+}@C_{74}^{3-}$  is different from divalent  $M^{2+}@C_{74}^{2-}$  EMFs because of the different electronic structure. Two questions are important in the study of the  $La@C_{74}$  system. Does the trivalent La ion show dynamic motion in the  $C_{74}$  ( $D_{3h}$ ) cage? And, what is the trajectory of the La ion in  $C_{74}$  ( $D_{3h}$ ) cage?

In this study, the interaction between the La atom and the  $C_{74}$  ( $D_{3h}$ ) cage is analyzed using all-electron relativistic DFT considering the representative patch of  $C_{74}$  ( $D_{3h}$ ). We optimized all possible isomers, searched for all transition states, and calculated the activation energies and the linear synchronous transit (LST). With the aid of the potential energy surface (PES) obtained by a single point energy scan over the  $\sigma_h$  surface of the cage, the motion of La in the cage is clarified. Furthermore, the trajectory of the La atom in the  $C_{74}$  ( $D_{3h}$ ) cage and the dynamic NMR spectra of  $La@C_{74}$  were obtained.

### Computational details

The DFT calculations of  $La@C_{74}$  were performed using the Dmol3 code [24, 25], and the Becke-Lee-Yang-Parr (BLYP) exchange correlation functional, which is a combination of the Becke exchange functional [26] coupled with the Lee-Yang-Parr (LYP) correlation potential [27]. The basis sets used in this work were double-numerical quality basis sets with polarization functions (DNP), which is comparable to the Gaussian 6–31 G\*\* basis set. It is known that relativistic effects play an important role in the chemical and physical properties of molecules containing heavy elements such as Lanthanide [28–33]. To take into account relativistic effects, the all-electron scalar relativistic method using the Douglas-Kroll-Hess (DKH) Hamiltonian, the most accurate approach available in Dmol3, was chosen. Self-consistent field procedures were carried out with a

convergence criterion of  $10^{-6}$  a.u. on the energy and electron density. Full geometry optimizations were performed using the Broyden-Fletcher-Goldfarb-Shanno (BFGS) algorithm [34] with a convergence criterion of  $10^{-3}$  a.u. on the displacement and  $10^{-5}$  a.u. on the energy. The transition states were searched using the complete LST/QST method with a convergence criterion of 0.002 a.u. on the RMS [35].

The NMR spectra of  $[La@C_{74}]^-$  was calculated using the ADF2008.01 program [36, 37], with the BLYP exchange correlation functional and the TZP basis set. The configurations were calculated by using the zero-order regular approximation (ZORA) [38–40], in which relativistic effects are considered.

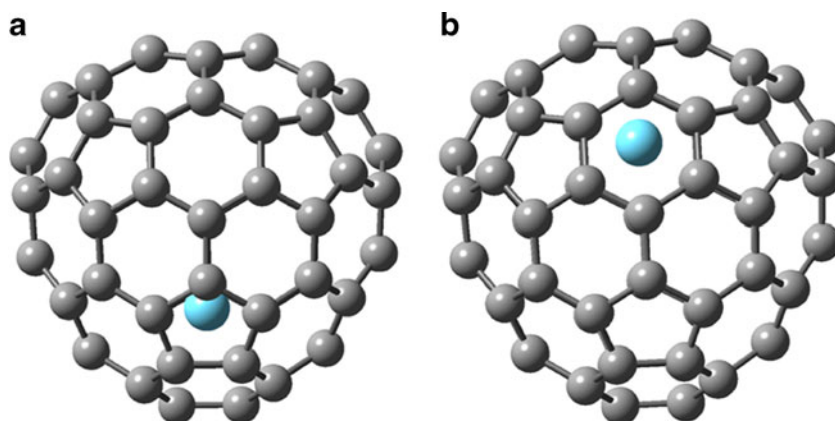
### Results and discussion

#### Geometry optimization and transition state search

The hollow  $C_{74}$  cage with  $D_{3h}$  symmetry was optimized, and then full geometry optimizations were carried out for  $La@C_{74}$  to find the favorable position for the La, with the most stable structure of  $La@C_{74}$  shown in Fig. 1a. The point group symmetry of  $La@C_{74}$  becomes  $C_{2v}$  upon encapsulation of the La atom. The La atom is located about 1.7 Å off-center and under a [6] double bond along a  $C_2$  axis on the  $\sigma_h$  plane. The vibrational frequencies of  $La@C_{74}$  were computed and all turned out to be real, ensuring that the configuration was a true minimum. Three symmetrically equivalent stable sites for the La atom were found inside the cage along the three  $C_2$  axes on the  $\sigma_h$  plane due to the  $D_{3h}$  symmetry of the cage.

Because there are three equivalent sites for the La atom, which are linked by rotation around the threefold axis, there is a transition state (TS) between two such equivalent minima. The TS structure (Fig. 1b) shows that La is located about 1.5 Å off-center along the  $C_2$  axis, and is under a hexagon opposite to the C=C double bond on the  $\sigma_h$  plane.

**Fig. 1** **a** Structure of  $La@C_{74}$  optimized by all-electron relativistic density functional theory Becke-Lee-Yang-Parr with double-numerical quality basis sets with polarization functions (DFT BLYP/DNP). **b** Optimized transition state



**Table 1** Calculated geometric and energetic parameters of La@C<sub>74</sub>-TS Transition state

La@C <sub>74</sub>	$R_{\text{La-C}}(\text{\AA})^a$	$E$ (kcal mol <sup>-1</sup> ) <sup>b</sup>	$E$ (kcal mol <sup>-1</sup> ) <sup>c</sup>	AE (kcal mol <sup>-1</sup> ) <sup>d</sup>
Energy stationary point	2.518	0	278.1	0
TS	2.585	19.3	277.1	18.4

<sup>a</sup> Shortest La-C distance<sup>b</sup> Relative total energy<sup>c</sup> Zero-point vibration energy<sup>d</sup> Activation energy

The structure of the transition state also has  $C_{2v}$  symmetry. The TS is a saddle point, since it has an imaginary frequency of  $63.5 \text{ i cm}^{-1}$ .

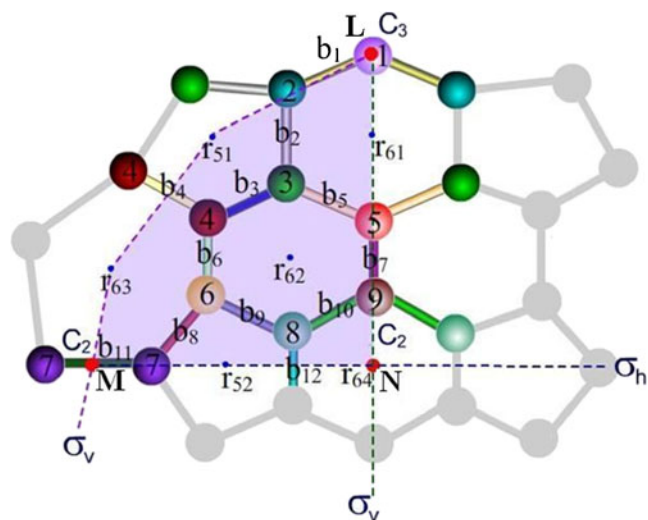
Table 1 shows the shortest La–C distances and the energy parameters of the two configurations. The lowest energy of the most stable configuration was taken as the relative zero point energy. The activation energy between the minimum energy point and the saddle point is  $18.4 \text{ kcal mol}^{-1}$ .

#### Interaction between La and C<sub>74</sub> ( $D_{3h}$ )

##### Representative patch for C<sub>74</sub> ( $D_{3h}$ )

C<sub>74</sub> fullerene has only one isomer that satisfies the isolated pentagon rule with  $D_{3h}$  symmetry. A right angle patch under a circum-spherical surface was taken as a representative patch of the C<sub>74</sub> ( $D_{3h}$ ) cage to describe the symmetry of a C<sub>74</sub> ( $D_{3h}$ ) cage simply as shown in Fig. 2.

The area of the patch LMN is equal to 1/12 of the total surface area. There are 27 key points on the C<sub>74</sub>( $D_{3h}$ ) representative patch. Points C1–C9 represent nine different

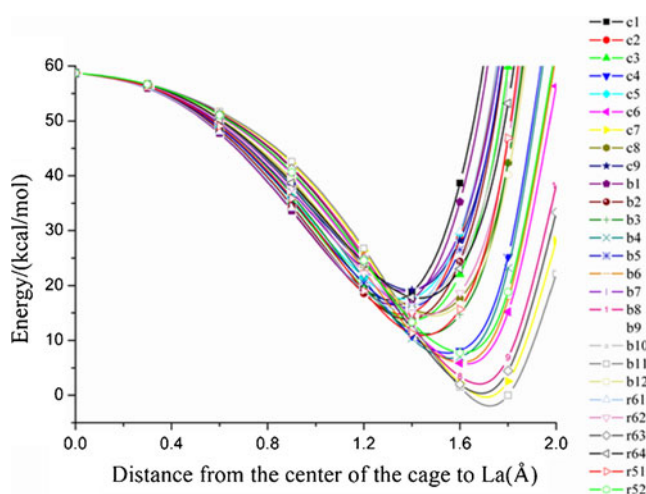
**Fig. 2** Representative patch of the C<sub>74</sub> ( $D_{3h}$ ) cage (shaded)

types of carbon atoms: C1 has  $C_{3v}$  local symmetry; C2, C5, C7, C9 have  $C_s$  local symmetry; and C3, C4, C6, C8 have  $C_1$  local symmetry. Points  $b_1$  to  $b_{12}$  represent the 12 distinct C–C bonds:  $b_{11}$  has  $C_{2v}$  local symmetry; and  $b_n$  ( $n=2, 3, 5, 6, 8, 9$ , and  $10$ ) have  $C_1$  local symmetry.  $r_{61}$  to  $r_{64}$  denote the four types of six-membered rings:  $r_{61}$  and  $r_{63}$  have  $C_s$  local symmetry;  $r_{64}$  has  $C_{2v}$  local symmetry.  $r_{51}$  and  $r_{52}$  represent the two types of five-membered rings with  $C_s$  local symmetry. When 12 elements of the  $D_{3h}$  group are operated on the patch, the patch will encompass the entire surface of the polyhedron.

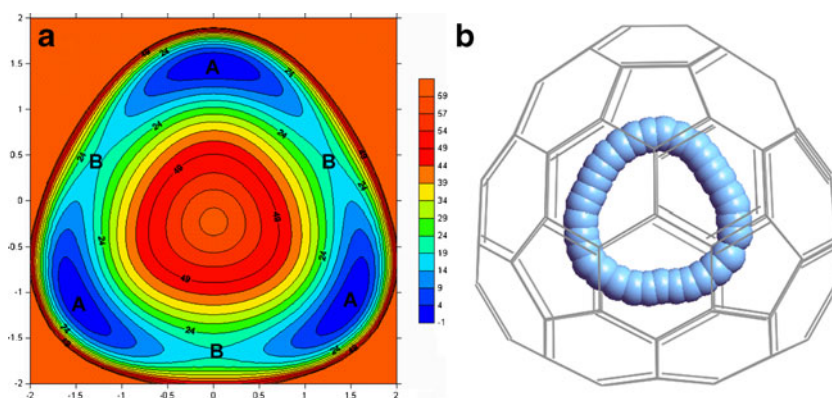
#### Interaction between La and C<sub>74</sub> ( $D_{3h}$ )

The interaction energy of La–C<sub>74</sub>( $D_{3h}$ ) was calculated. Firstly, the C<sub>74</sub>( $D_{3h}$ ) cage was fixed, and then the La atom was allowed to approach the 27 key points along the radial directions, which pass through the center of the cage and the key points. The calculated potential energy curves as a function of the distance between the La atom and the key points are shown in Fig. 3. There is a minimum energy point in each potential energy curve. These energy minima are located between 1.6 and 1.8 Å from the center of the cage. The five minimum energy lines are C7,  $b_{11}$ ,  $r_{52}$ ,  $b_{12}$  and  $r_{64}$ , all of which are located in the  $\sigma_h$  plane. To clarify the minimum energy pathway, the  $\sigma_h$  plane should be studied in detail.

The PES of the  $\sigma_h$  plane was obtained using the single point energy scan as shown in Fig. 4a. There are three equivalent local minima on the  $\sigma_h$  surface, denoted by A, that are about 1.7 Å off-center. The transition states B are about 1.5 Å off-center. As shown in Fig. 4a, the relative energy between A and B is about  $18.4 \text{ kcal mol}^{-1}$ , which is consistent with the value from the geometry optimization listed in Table 1. The lowest energy pathway is in the  $\sigma_h$

**Fig. 3** Calculated potential energy curves of La atoms approaching the 27 key points of the C<sub>74</sub> cage (BLYP/DNP)

**Fig. 4** **a** Calculated potential energy surface (PES) of La atom on the  $\sigma_h$  surface of the  $C_{74}$  ( $D_{3h}$ ) cage at BLYP/DNP level (energy in eV). **b** Probable trajectory of the La atom in the cage

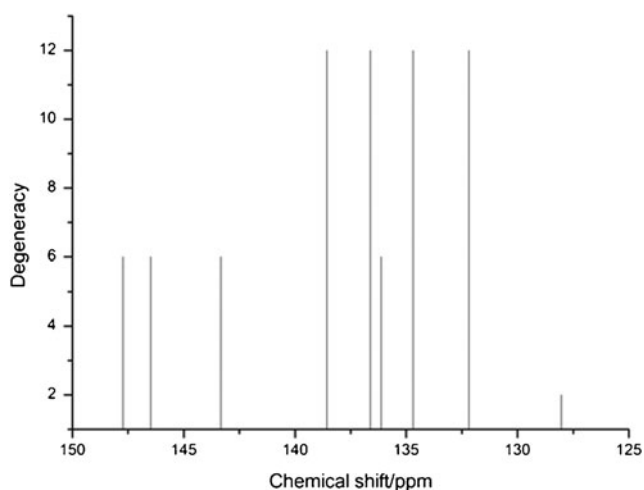


plane(A-B-A-B-A-B-A). It is proposed that the motion of La forms a ring. The motion of La in the  $C_{74}$  ( $D_{3h}$ ) cage forms a ring as shown in Fig. 4b. The trajectory of La has  $D_{3h}$  symmetry and is in the  $\sigma_h$  plane. It is reasonable to predict the motion of the atom inside the cage by comparing the energy of different configurations.

#### Prediction of $^{13}\text{C}$ NMR spectrum

Using the trajectory of La in the  $C_{74}$  ( $D_{3h}$ ) cage, the  $^{13}\text{C}$  NMR spectrum of  $\text{La}@C_{74}$  can be predicted. First, we qualitatively analyze the  $^{13}\text{C}$  NMR spectrum of  $\text{La}@C_{74}$ . There should be nine lines that correspond to the nine types of atoms in the  $C_{74}$  ( $D_{3h}$ ) cage. The nine carbon atoms are defined as shown in Fig. 2. We can predict the intensity of lines: C1 has  $C_{3v}$  local symmetry, so it should give 1/6 intensity signals; C2, C5, C7, C9 have  $C_s$  local symmetry, so should give half intensity signals; C3, C4, C6, C8 have  $C_1$  local symmetry, so should give full intensity signals.

In addition, we calculated the average chemical shifts for all of the carbon atoms to simulate the dynamic  $^{13}\text{C}$  NMR



**Fig. 5** Calculated  $^{13}\text{C}$  NMR spectrum of  $[\text{La}@C_{74}]^-$

spectrum of  $\text{La}@C_{74}$ . The 74 C atoms were divided into nine groups according to the representative patch of  $C_{74}$  ( $D_{3h}$ ), and their arithmetic average chemical shifts were obtained. Figure 5 shows the  $^{13}\text{C}$  NMR spectrum of  $[\text{La}@C_{74}]^-$ , with the calculated chemical shifts ranging from 128 to 147.7 ppm. The calculated chemical shieldings  $[\delta(\text{Ci})]$  were then referenced to those of  $C_{60}$  to obtain the calculated NMR chemical shifts using the following formula [41]:

$$\delta(\text{Ci}) = \delta(\text{C}_{60}) + \sigma(\text{C}_{60}) - \sigma(\text{Ci}), \quad (1)$$

where  $\delta(\text{C}_{60})$  is taken to be 143.15 ppm [42],  $\sigma(\text{C}_{60})$  is the calculated chemical shielding of  $C_{60}$  (28.30 ppm), and (Ci) is the carbon atom under consideration.

From Fig. 5, the  $^{13}\text{C}$  NMR spectrum of  $[\text{La}@C_{74}]^-$  contains nine signals. The spectrum has four lines with full intensity, four lines with half intensity, and one line with 1/6 intensity. This means that there are nine types of carbon atoms in the metallofullerene cage. This NMR pattern indicates that  $\text{La}@C_{74}$  retains the  $D_{3h}$  symmetry of the  $C_{74}$  cage. These spectral features are similar to those of  $\text{Ca}@C_{74}$  [22] and  $\text{Yb}@C_{74}$  [4].

Computed [43] and experimental electronic, infrared and Raman spectra indicates that  $\text{Eu}@C_{74}$  has a  $C_{2v}$  symmetric structure derived from the  $D_{3h}$  symmetric IPR isomer of  $C_{74}$ . The results of a study by Rappoport and Furche [43] together with those of the present study demonstrate that computational works can help to interpret spectroscopic data and to assign structures.

#### Conclusions

In summary, we have investigated  $\text{La}@C_{74}$  by all-electron relativistic DFT. We obtained the lowest energy structure of  $\text{La}@C_{74}$  and the TS. Both configurations of the lowest energy structure and the TS have  $C_{2v}$  symmetry. The energy barrier is  $18.4 \text{ kcal mol}^{-1}$  for La to hop from one stable site to another. The lowest energy pathway is in the  $\sigma_h$  plane,



and the trajectory is a ring with  $D_{3h}$  symmetry. On the basis of this, the dynamic  $^{13}\text{C}$  NMR spectrum of  $\text{La}@C_{74}$  was predicted. There are 9 lines in the spectrum: four of them have full intensity, four have half intensity, and one has one-sixth intensity. This NMR pattern indicates that the system has  $D_{3h}$  symmetry, which is the same as the surrounding cage. This indicates the dynamic motion of the encaged metal atom and can be used to determine the dynamic motion of one or more atoms inside carbon cages.

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